# Some Notes on the Skwirzynski Transformation

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These notes are intended to be an extension to *Chapter 2.8.2 The 'Skwirzynski' transformation*, by R. Nouta. In here, the importance of the Skwirzynski transformation when designing realizable even order ladder filter structures with lumped elements has been described.

The description of the transfer function of an even order low pass filter –resulting from the theory– is commonly denoted as a type A design. Realizable filters will be of Type B or Type C. Type B corresponds to Nouta's **Example 3, Possibility 1**, while Type C is described in **Example 3, Possibility 2**. Skwirzynski denotes a Type B transformation as the 'wiggle' transformation, indicated with  $\Omega \to \widetilde{\Omega}$ . In his terminology, a Type C transformation is a 'bar' transformation:  $\Omega \to \overline{\Omega}$ .

#### Type A

For illustration, a stylized transfer curve of a 6th order normalized Cauer filter is used with the characteristic frequency points denoted in Figure 1. Note the specific numbering of the  $\Omega_p$  versus the  $\Omega_s$  frequency points, which indicate the relationship between equally numbered points.





In the literature we encounter two frequently used, different choices for normalization

1. method 1, for reasons of symmetry (used a.o. by Nouta, Skwirzynski and Antoniuou):

$$\Omega_P = \sqrt{k}, \quad \Omega_C = 1 \text{ and } \Omega_S = \frac{1}{\sqrt{k}} \quad (k < 1)$$

2. method 2, for reasons of comparability with the Chebyshev normalization (used by Christian & Eisenmann):  $\Omega_P = 1, \quad \Omega_C = \frac{1}{\sqrt{k}} \text{ and } \Omega_S = \frac{1}{k} \quad (k < 1)$ 

## Type B

A new frequency variable  $\omega$  will be used that is obtained by shifting  $\Omega_{s1}$  to infinity ( $\omega_{p1} = \infty$ ), while keeping the frequency points at 0 and 1 in place as illustrated in Figure 2.



Figure 2.

Now the new frequency variable can be expressed in terms of the old one with

$$\omega = \Omega \sqrt{\frac{\Omega_{s1}^2 - 1}{\Omega_{s1}^2 - \Omega^2}}$$

If P are the complex poles of the transfer function, then the new poles p can be obtained with

$$p=P\sqrt{\frac{\Omega_{s1}^2-1}{\Omega_{s1}^2+P^2}}$$

Since usually  $\Omega_{s1}$  has been derived from  $\Omega_{p1}$ , it will be more accurate to use just this  $\Omega_{p1}$ . Considering the two normalization methods, we can write

$$\omega = \Omega \sqrt{\frac{1 - C_k \Omega_{p1}^2}{1 - C_k \Omega_{p1}^2 \Omega^2}} \quad \text{and} \quad p = P \sqrt{\frac{1 - C_k \Omega_{p1}^2}{1 + C_k \Omega_{p1}^2 P^2}}$$
where 
$$\frac{\text{normalization}}{C_k | 1 | k^2}$$

since in method 1  $\Omega_{s1} = \frac{1}{\Omega_{p1}}$ , while in method 2  $\Omega_{s1} = \frac{1}{k\Omega_{p1}}$ .

## Type C



Except for the shift of  $\Omega_{s1}$  to infinity,  $\Omega_{p1}$  will be shifted to 0 as well.

Figure 3.

The new frequency variable becomes

$$\omega = \sqrt{\frac{\left(\Omega_{s1}^2 - 1\right)\left(\Omega^2 - \Omega_{p1}^2\right)}{\left(1 - \Omega_{p1}^2\right)\left(\Omega_{s1}^2 - \Omega^2\right)}}$$

and the new locations of the poles become

$$p = -\sqrt{\frac{\left(\Omega_{s1}^2 - 1\right)\left(P^2 + \Omega_{p1}^2\right)}{\left(1 - \Omega_{p1}^2\right)\left(\Omega_{s1}^2 + P^2\right)}}$$

Again realizing that  $\Omega_{s1}$  has in fact been derived from  $\Omega_{p1}$ , we will eliminate  $\Omega_{s1}$ . With the different normalization methods in mind, we write

$$\begin{aligned} \omega &= C_m \sqrt{\frac{\Omega^2 - \Omega_{p1}^2}{1 - C_k \Omega_{p1}^2 \Omega^2}} \quad \text{and} \quad p = -C_m \sqrt{\frac{P^2 + \Omega_{p1}^2}{1 + C_k \Omega_{p1}^2 P^2}} \\ \end{aligned} \\ \text{where} \quad \frac{\frac{\text{normalization}}{\text{method 1} \quad \text{method 2}}}{\frac{C_m}{C_k} \quad 1 \quad \sqrt{\frac{1 - k^2 \Omega_{p1}^2}{1 - \Omega_{p1}^2}}} \\ \frac{C_m + 1}{C_k + 1 \quad k^2} \end{aligned}$$

The equations above have been implemented and tested in several MATLAB programs, viz. d:\filters\cauer\skwirBC.m (first test program), d:\filters\cauer\data6ABC.m (comparison of the output data of 6th order filters), and the final function d:\filters\cauer\Cauer.m with syntax [Ws,wp,ws,p] = Cauer(N,rp,rs,normth<,ftype>) where input parameters N : the order of the filter rp : the maximum ripple in the pass band in dB rs : the minimum attenuation in the stop band in dB normth : the normalization method, 0 for  $\Omega_C = 1$  or 1 for  $\Omega_P = 1$ ftype : for even order filters, type 'A', 'B' or 'C' needs to be specified output parameters  $\texttt{Ws} \ : \ \ \Omega_S \text{ or } \omega_S$ wp : the  $\Omega_p$  or  $\omega_p$  vector ws : the  $\Omega_s$  or  $\omega_s$  vector p : the complex poles of the Cauer function

Note: wp corresponds with the ATTEN.ZEROS, and ws with the ATTEN.POLES listed in Christian & Eisenmann

### **References:**

E. Christian & E. Eisenmann, Filter Design Tables and Graphs, John Wiley & Sons, Inc. 1966

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J.K. Skwirzynski, Design Theory and Data for Electrical Filters, Van Nostrand Company Ltd, 1965